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## Transient analysis of the slamming wave load on an offshore wind turbine foundation generated by different types of breaking waves

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### Abstract

In this study, a computational fluid dynamics model was employed to simulate the impact of a breaking wave on the foundation of an offshore wind turbine. The accuracy of the numerical simulation was validated in comparison with published experimental results for the Large Wave Channel (GWK) at Hannover, Germany. To understand the influence of the wave-breaking location on the monopile support structure, a series of shoaling waves with different spilling and plunging breaking wave types were considered in terms of the surf-similarity parameter ( $\xi_0$ )<sup>4</sup> and relative breaking distance ( $x'_b$ ). After processing the total wave load with an empirical mode decomposition method and a low-pass filter, the instantaneous nature of the slamming wave load generated by different breaking wave types was obtained. Subsequently, an exponential-type refined dynamic model was established to consider the characteristics of the filtered slamming wave load.  $\xi_0$  and  $x'_b$  were found to be dominant factors for the run-up phenomena, slamming characteristics, and wave-dependent parameters. The results demonstrated that these two coefficients can be applied to determine the nature of the run-up and slamming wave loads.

### Wave Conditions

The surf similarity parameter  $\xi_0$  is used to determine the wave-breaking type. Twenty Iribarren numbers in the ranges of  $\xi_0 < 0.5$  (Spilling) and  $0.5 < \xi_0 < 2.3$  (Plunging) are presented in Fig.1(a-b) and Fig.1(c-d), respectively. These wave conditions are simulated under a water depth  $d$  of 3.8 m, sloping bottom angle  $\beta$  of 5.71°, and cylinder axis position of 45 m from the inlet.

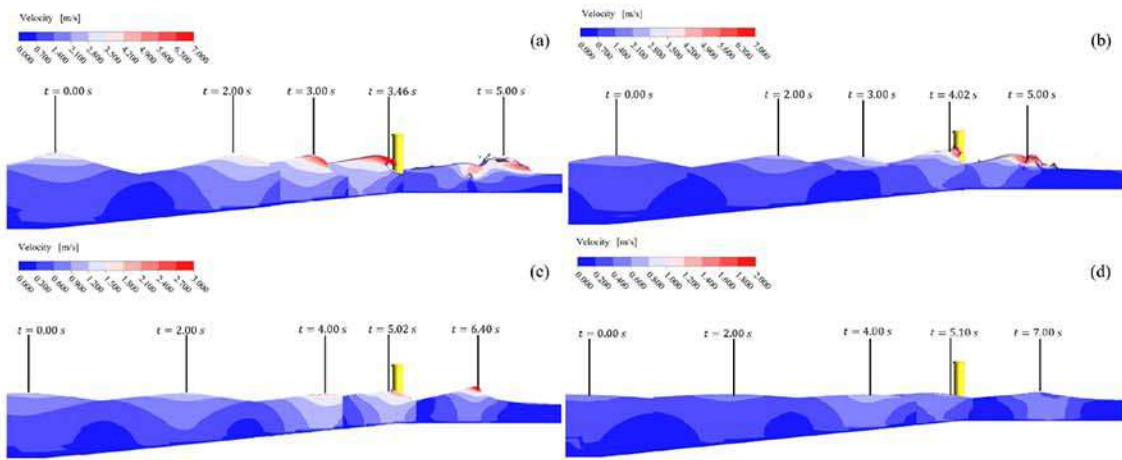


Fig.1 Scenarios of waves propagating over the sloping bottom for cases  $\xi_0 =$  (a) 0.35, (b) 0.41, (c) 0.51 and (d) 0.72.

### Numerical Setup

The grid resolution is a crucial factor influencing the accuracy of the calculations obtained using the VOF method<sup>2</sup>. The mesh is generated by refined to 0.05 m to capture the surface elevation [Fig.2 (a)]. A numerical beach treatment is conducted in front of the outlet boundary to reduce the numerical reflection of waves [Fig. 2 (b)].

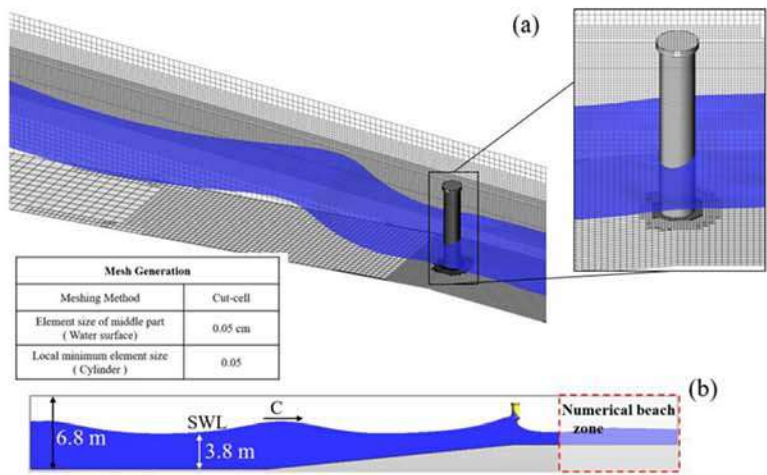


Fig.2 (a) Mesh generation; and (b) numerical beach treatment.

### Result and Discussion

To test the sensitivity of the CFD model to the grid size and time step, the wave conditions of  $H_0 = 1.3$  m and  $T_0 = 4$  s were used to evaluate the wave height at the wave-breaking point. A series of convergence tests with different combinations of grid sizes and time steps were performed with simultaneous refinement<sup>6</sup>.

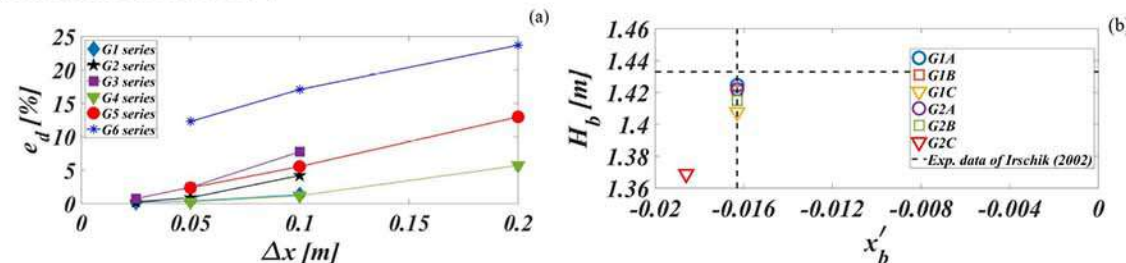


Fig. 3. (a) Spatial and temporal convergences for the wave breaking height, and (b) comparisons of the simultaneous grid and time refinement approach and the experimental data of Irschik (2002)<sup>5</sup>.

The properties of breaking wave types for regular waves with the surf-similarity parameters  $\xi_0 = 0.35, 0.41, 0.51,$  and  $0.72$  are investigated to identify different scenarios and locations of wave breaking, as shown in Fig. 1(a)-(d). Fig. 4(a)-(d) present the refined slamming wave load model<sup>9</sup> with a normalized time scale for the cases of  $H_0 = 1.8$  m, 1.3 m, 0.8 m, and 0.4 m, respectively. The fitting curves were obtained by linear regression with a coefficient of determination<sup>7,8</sup> (i.e.,  $R^2$  value) of 0.95. In this case, both the exponential growth and decay parameters (i.e.,  $\alpha_1$  and  $\alpha_2$ ) can be estimated with a high confidence level.

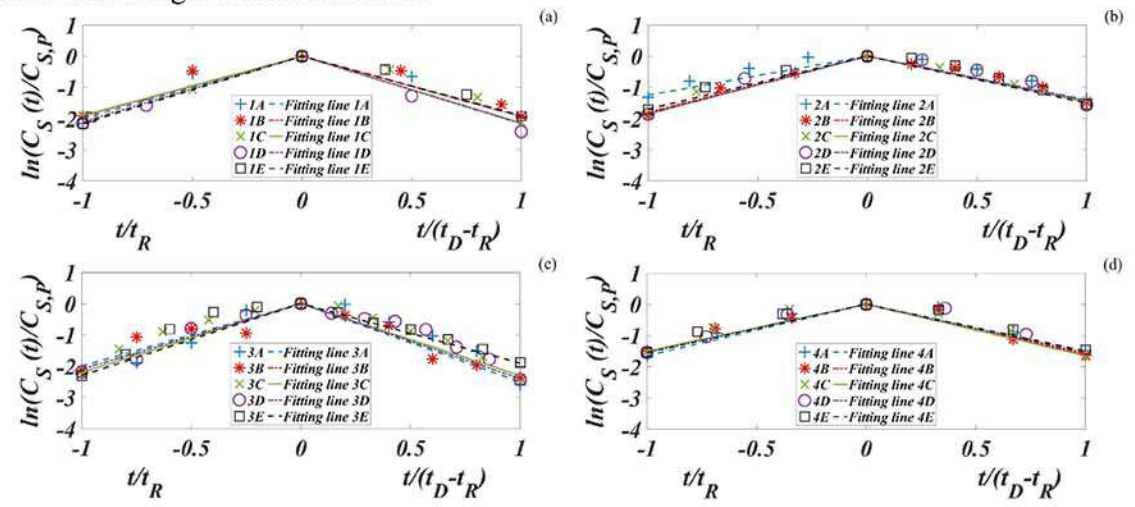


Fig.4 Refined slamming wave load models for Cases (a)  $H_0=1.8$  m, (b)  $H_0=1.3$  m, (c)  $H_0=0.8$  m, and (d)  $H_0=0.4$  m.

By mean of our CFD model, the corresponding  $C_{S,P}$  and  $\lambda$  values with respect to  $x'_b$  are shown in Fig. 5 (a) and (b), respectively. Furthermore, wave-dependent parameters<sup>9</sup> (i.e.,  $\omega_1, \omega_2,$  and  $\omega_3$ ) are presented in Fig. 5 (c), (d) and (e), respectively. Fig. 5 (f) show the normalized maximum run-up height  $R_{u,max}/\eta_{max}$ .

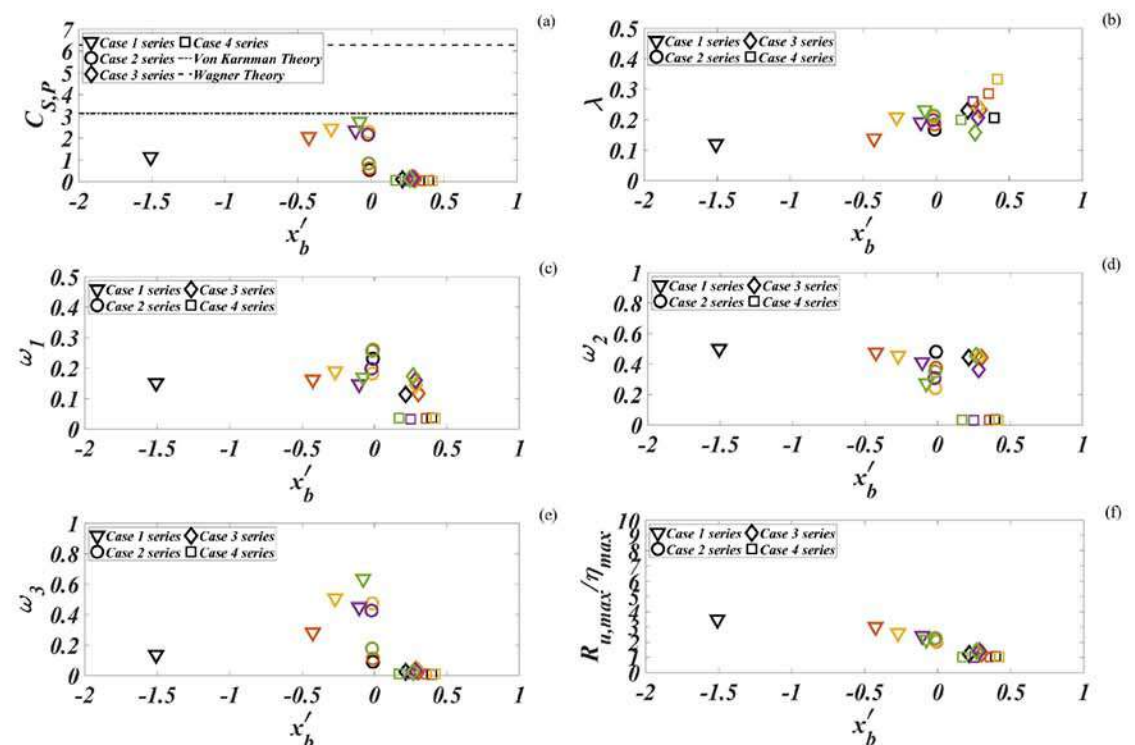


Fig.5 Variations of (a)  $C_{S,P}$ , (b)  $\lambda$  (c)  $\omega_1$ , (d)  $\omega_2$ , (e)  $\omega_3$ , and (f)  $R_{u,max}/\eta_{max}$  with respect to  $x'_b$

### Conclusions

1. The numerical setting of  $\Delta x = 0.025$  m and  $\Delta t = 0.0025$  s (i.e.,  $C_r = 0.5$ ) was in good agreement with the experimental data of Irschik<sup>56</sup> in terms of the form, load, and breaking location of the wave interacting with the monopile foundation in the wave flume.
2. According to the simulation results, the maximum slamming coefficient  $C_{S,P}$  obtained from spilling breaking waves was found to approach Von Karman's value rather than Wagner's value<sup>3</sup>. The relative breaking distance  $x'_b$  has a considerable influence on  $C_{S,P}$  and  $R_{u,max}/\eta_{max}$ , which have opposite trends when  $x'_b < 0$ . Increasing  $x'_b$  caused  $\lambda$  to increase and  $R_{u,max}/\eta_{max}$  to decrease. The impact parameter  $\omega_3$  showed a significant response in the range of  $-0.5 < x'_b < 0$  but decreased rapidly to a negligible value when  $x'_b > 0$ .  $\omega_3$  is dominated by  $C_{S,P}$  instead of  $\lambda$ .

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